# Data Extraction with ${ }^{3} \mathrm{He}$ detectors 

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## Introduction

In [1] and [2] we have looked at the performance of the SNO detector with the 'default' neutral current detection option ( NaCl dissolved in the $\mathrm{D}_{2} \mathrm{O}$ ). Three different $\Delta m^{2}, \sin ^{2} 2 \theta$ scenarios and the case of vacuum oscillations were investigated. Data were generated to simulate a 3 year experiment : a light water fill, a heavy water fill and finally a heavy water plus NaCl fill.

In this paper we will compare a similar 3 year experiment for a SNO detector which uses ${ }^{3} \mathrm{He}$ proportional counters to detect the neutrons created by the Neutral Current reaction.

## Calculations

The standard Monte Carlo code [1-5] was modified to include 109 cylindrical tubes mounted vertically on a 99 cm spacing grid in the $\mathrm{D}_{2} \mathrm{O}$. The tubes were assumed to have an outside diameter of 5 cm , to have a refractive index of 4.0 and to be highly absorbing to photons with a wavelength between

300 and 700 nm . A material with a refractive index of 4.0 would reflect $25 \%$ at normal incidence when immersed in water. The present design calls for thin acrylic tubes coated with a metal on the inside. Such tubes would likely have a lower refractive index so the calculation probably represents a 'worst case scenario'.

## Energy calibration/resolution and reconstruction effects

Two questions immediately come to mind when one introduces the idea of ${ }^{3} \mathrm{He}$ counters in the $\mathrm{D}_{2} \mathrm{O}$ : how much light is lost and will the tubes destroy the spherical symmetry ? Figure la shows the reconstructed energy for ten thousand 10.511 MeV electrons started uniformly throughout the $\mathrm{D}_{2} \mathrm{O}$ for the standard SNO detector. Figure $1 b$ shows the same with the ${ }^{3} \mathrm{He}$ counters present. We clearly lose both light and resolution. The number of hits per MeV goes from 9.5 to 8.1 but the width of the distribution stays the same, meaning that the percentage resolution gets worse.

To test directional effects simulations were done for the ${ }^{3} \mathrm{He}$ tube configuration with 10.511 MeV electrons started near the center of the $\mathrm{D}_{2} \mathrm{O}$ in three different directions: along the $z$-axis, in the x - y plane $45^{\circ}$ to the x -axis and along the x -axis (figures 2 to 4 ). There is little difference in either the cosine of the angle between the electron and the reconstructed direction and in the reconstructed energy (table 1).

## Backgrounds

As in [1] we use the 'White Book' [6] radioactivity levels for the detector components (table 2) with the added assumption that the major radioactivity in the ${ }^{3} \mathrm{He}$ counters is the acrylic tubing ( $1 / 8$ inch wall thickness). If we take the 'White Book' numbers for acrylic then the 109 counter contribute a little more than $10 \%$ of the $\mathrm{D}_{2} \mathrm{O}$ in terms of neutrons. This is comparable to the contribution from the NaCl in the default option. For simplicity we have therefore used the same neutron production numbers
for the ${ }^{3} \mathrm{He}$ tubes as were used for the NaCl in [1]. Based on these and a production rate from the sun of 4500 neutrons per year [1] we arrive at table 3.

## Extraction of data

Figure 5 shows a summary of signals and backgrounds (per year) reconstructed inside a radius of 600 cm assuming vacuum oscillations and a total flux of $5.8 \times 10^{6} \mathrm{~cm}^{2} \mathrm{~s}^{-1}$. This is to be compared with figure 1 in [1] (reproduced here as figure 6). The threshold is somewhat higher but the smaller Neutral Current (Čerenkov) signal helps in sorting out the different components.

For the analysis we follow the same procedure as in [1] where the data was sorted according to angle and energy ( $40 \times 200$ bins ) and the angular projections were fitted with the elastic scattering angular distribution plus a term of form $1-\frac{1}{3} \times \cos \theta$ plus a constant (representing the sum of the Neutral Current part and the background). In contrast with [1] year 2 and 3 has the detector in the same configuration and consequently we just sum the two. The analysis is done for an outer radius of 600 cm only.

Some typical fits to the sum of years 2 and 3 are shown in figures 7 b to 7 e as a function of lower threshold. Figure 7 a is the corresponding energy projection for all values of the angle. At 40 hits we include a sizable background which decreases very fast with increasing threshold. However at high $\mathrm{N}_{\text {hit }}$ there is not much of an ES signal left. Fits for a fixed threshold of $N_{h i t}=50$ are displayed in figures 8 to 11 for vacuum oscillations, $\Delta \mathrm{m}^{2}$ $=10^{-6}, 10^{-5}$ and $10^{-4}$ for year 1 and years $2+3$.

The fits are summarized in tables 4 and 5 where we list the lower threshold, the fitted values and the Monte Carlo input numbers plus the total $\chi^{2}$. The same numbers are presented in ratio form in table 6 . At around $\mathrm{N}_{\text {hit }}=50$ we more or less extract what is put in.

Expected rates per year, based on the high statistics multiyear runs, are listed in table 7 where we also include the expected number from the
${ }^{3} \mathrm{He}$ counters. We assume an efficiency of $85 \%$ for these.

Ratios of extracted signals to that expected from the standard solar model are summarized in table 8 . The NC fux listed is that determined from the capture in the deuterium not the one determined from the ${ }^{3} \mathrm{He}$ counters. Due to the poor statistics and the high correlation with the CC spectrum these numbers are virtually meaningless. The other ratios track very well and if the uncertainty in doing the subtraction of the neutron background in the ${ }^{3} \mathrm{He}$ counters is low enough (overall uncertainty less than about $15 \%$ ) then we could prove the presence of oscillations for any of the 4 cases studied at the $3 \sigma$ level.

Another way of analyzing the data is to form the ratio of combinations of sums and differences of the total number of counts above some threshold similarly to what was done in [1]. This assumes constant fluxes and constant backgrounds as a function of time. If $T_{i}$ is the total number of counts in year $i$ and $H$ is the average number of counts in the ${ }^{3} \mathrm{He}$ counters then the ratio $r=\frac{H}{T_{3}-T_{2}-2 \times T_{1}}$ is related to $R^{\prime}=H_{o b s} / C C_{o b s}$ by $R^{\prime}=\frac{r}{1-\pi \times r}$ where $n$ is the ratio of the neutron detection efficiencies for deuterium and ${ }^{3} \mathrm{He}$. It can be determined accurately with a source.

Figure 12 shows

$$
R=\left[H_{o b s} / C C_{o b s}\right] \times\left[C C_{m c} / H_{m c}\right]
$$

where $C C_{m c}$ is evaluated based on the no oscillation scenario for a lower threshold on $N_{\text {hit }}=50$ plotted as function of the neutron background in units of 965 detected in the ${ }^{3} \mathrm{He}$ counters per year. This is the 'White Book' prediction; however it does not include any $\alpha$ background in the ${ }^{3} \mathrm{He}$ detectors. As in [1] we assume a $20 \%$ uncertainty in determining the total neutron background. The uncertainty in $H_{o b s}$ is just the statistical one. There is a small improvement over the equivalent figure for NaCl figure 13 (figure 11 from [1]). In both cases we have assumed fluxes which do not vary as a function of time. The reason why we get approximately the same answer is that the correlation between the charged current and the neutral current components is essentially zero when one does subtractions. Simultaneous fit of CC and NC on the other hand introduces a large correlation coefficien:- which increases the uncertainty significantly.

## Conclusions

We have shown that the ${ }^{3} \mathrm{He}$ tubes do not interfere significantly with the Čerenkor light and that it is still possible to get a reasonable signal from the CC and ES reactions. The amplitude of the ES signal can be extracted over a wide range of MSW parameters from the angular fits alone as can the CC signals. In terms of answering the question of whether any of the four chosen scenarios differ from the no oscillation case by more than $3 \sigma$ the ${ }^{3} \mathrm{He}$ option is doing a little better than NaCl with straight subtraction. In both cases we have assumed that the total neutron background could be determined by other means to $20 \%$ but we have not specifically included $\alpha$-backgrounds or any other systematic uncertainties in the ${ }^{3} \mathrm{He}$ counters.

## References

[1] SNO-STR-91-06
[2] SNO-STR-91-28
[3] SNO-STR-89-01
[4] SNO-STR-89-43
[5] SNO-89-15
[6] SNO-87-12

| $\mathrm{e}^{-}$ <br> direction | Reconstructed <br> $\mathbf{N}_{\text {hit }}$ | Reconstructed <br> $\cos$ |
| :---: | :---: | :---: |
| $(0,0,1)$ | $87.3(11.4)$ | $.91(.11)$ |
| $\left(\frac{1}{\sqrt{2}} \cdot \frac{1}{\sqrt{2}}, 0\right)$ | $84.1(11.1)$ | $.89(.12)$ |
| $(1,0,0)$ | $84.0(11.1)$ | $.89(.11)$ |

Table 1: Average number of reconstructed hits and average cosine of angle between $\mathrm{e}^{-}$direction and reconstructed direction. Numbers in brackets are second moments. A total of 1,000 electron showers ( 10.511 MeV total energy) were started at $(50,0,0)$.

|  | Mass(tonnes) | $\mathrm{Th}(\mathrm{g} / \mathrm{g})$ | $\mathrm{U}(\mathrm{g} / \mathrm{g})$ |
| :--- | :---: | :---: | :---: |
|  |  |  |  |
| $\mathrm{D}_{2} \mathrm{O}$ | 1000 | $11 \times 10^{-15}$ | $11 \times 10^{-15}$ |
| Acrylic | 30.0 | $1.9 \times 10^{-12}$ | $3.6 \times 10^{-12}$ |
| $\mathrm{H}_{2} \mathrm{O}$ | 1667.7 | $22 \times 10^{-15}$ | $15 \times 10^{-15}$ |
| PMT | 7.5 | $0.1 \times 10^{-6}$ | $0.1 \times 10^{-6}$ |

Table 2: Masses and levels of ${ }^{232} \mathrm{Th}$ and ${ }^{238} \mathrm{U}$ in various detector components

| Component | Production rate <br> per year | Captures per year <br> ${ }^{3} \mathrm{He}$ |  |
| :---: | :---: | :---: | :---: |
| deuterium |  |  |  |
| Sun | 4500 | 2049 | 566 |
| Int. background | 1432 | 652 | 180 |
| Ext. background | 2535 | 483 | 127 |

Table 3: Neutron production rates and capture rates per year. Internal background refers to neutrons arising from the $\mathrm{D}_{2} \mathrm{O}$ and the ${ }^{3} \mathrm{He}$ counters. The rest (acrylic vessel, $\mathrm{H}_{2} \mathrm{O}$, PMT's etc. is included in the external component

| $\begin{aligned} & \Delta \mathrm{m}^{2} l \\ & \sin ^{2} 2 \theta \end{aligned}$ | $\begin{aligned} & \text { Thres. } \\ & \mathrm{N}_{\text {hit }} \end{aligned}$ | ES |  | Background |  | $\begin{gathered} \chi^{2} \\ (40 \mathrm{pts} .) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Fit | M.C. | Fit | M.C. |  |
| vac. | 40 | 307 (25) | 336 | 1137 (38) | 1108 | 121. |
|  | 50 | 202 (15) | 217 | 62 (10) | 47 | 33.1 |
|  | 60 | 123 (11) | 123 | 5 (3) | 5 | 38.2 |
|  | 70 | 74 (9) | 74 | 2 (2) | 2 | 34.0 |
|  | 40 | 546 (30) | 575 | 1137 (38) | 1108 | 114. |
| $10^{-6}$ | 50 | 360 (20) | 375 | 62 (10) | 47 | 29.1 |
| 0.01 | 60 | 223 (15) | 22.5 | 8 (4) | 5 | 32.3 |
|  | 70. | 132 (12) | 133 | 3 (2) | 2 | 36.3 |
|  | 40 | 39 (16) | 53 | 1122 (37) | 1108 | 132. |
| $10^{-5}$ | 50 | 35 (7) | 35 | 47 (8) | 47 | 35.2 |
| 0.3 | 60 | 22 (5) | 21 | 4 (3) | 5 | 35.6 |
|  | 70 | 10 (3) | 10 | 2 (2) | 2 | 32.7 |
|  | 40 | 155 (20) | 175 | 1128 (37) | 1108 | 130. |
| $10^{-4}$ | 50 | 73 (10) | 82 | 56 (9) | 47 | 27.4 |
| 0.01 | 60 | 31 (6) | 31 | 54 (3) | 5 | 31.3 |
|  | 70 | 13 (4) | 13 | 3 (2) | 2 | 35.6 |

Table 4; Year 1 data for $0 \leq \mathrm{r} \leq 600 \mathrm{~cm}$. Angular distribution fitted to $\mathrm{ES}(\theta)+$ constant.

| $\begin{aligned} & \Delta \mathrm{m}^{2} / \\ & \sin ^{2} 2 \theta \end{aligned}$ | Thres. <br> $\mathrm{N}_{\text {hit }}$ | ES |  | CC |  | Background + NC |  | $\begin{gathered} \chi^{2} \\ (40 \mathrm{pts}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Fit | M.C. | Fit | M.C. | Fit | MC |  |
| vac. | 40 | 711 (59) | 614 | 6708 (659) | 58.53 | 2207 (692) | $2263+895$ | 59.7 |
|  | 50 | 499 (43) | 412 | 4.535 (475) | 4486 | 242 (497) | $119+259$ | 23.7 |
|  | 60 | 295 (32) | 251 | 2967 (374) | 3043 | 86 (389) | $12+42$ | 29.2 |
|  | 70 | 126 (14) | 141 | 1.500 (285) | 18.54 | 374 (295) | $3+2$ | 47.6 |
|  | 40 | 1104 (71) | 1016 | 9205 (764) | 8807 | 2673 (803) | $2263+895$ | 99.2 |
| $10^{-6}$ | 50 | 741 (52) | 647 | 7272 (581) | 681.5 | -173 (605) | $119+259$ | 54.4 |
| 0.01 | 60 | 439 (40) | 402 | 4789 (465) | 4693 | -79 (482) | $12+42$ | 48.0 |
|  | 70 | 278 (29) | 237 | 3476 (3.53) | 28.5 | -667 (363) | $3+2$ | 36.4 |
|  | 40 | 104 (37) | 118 | 1450 (439) | 10.3 | 2745 (464) | $2263+895$ | 157. |
| $10^{-5}$ | 50 | 83 (20) | 75 | 810 (229) | 767 | 327 (240) | $119+259$ | 36.4 |
| 0.3 | 60 | 40 (14) | 43 | 317 (159) | 506 | 247 (166) | $12+42$ | 39.0 |
|  | 70 | 20 (9) | 26 | 228 (113) | 285 | 68 (117) | $3+2$ | 42.8 |
|  | 40 | 324 (46) | 319 | 2922 (517) | 2246 | 2694 (545) | $2263+895$ | 105. |
| $10^{-4}$ | 50 | 147 (23) | 136 | 1066 (255) | 1005 | 308 (267) | $119+259$ | 37.9 |
| 0.01 | 60 | 70 (13) | 49 | 606 (143) | 389 | -182 (147) | $12+42$ | 39.7 |
|  | 70 | 16 (7) | 18 | 163 (80) | 136 | -18 (83) | $3+2$ | 41.1 |

Table 5: Year $2+3$ data for $0 \leq r \leq 600 \mathrm{~cm}$. Angular distribution fitted to $\mathrm{ES}(\theta)+\left[1-\frac{1}{3} \cos \right.$ $\theta]+$ constant.

| $\begin{gathered} \Delta \mathrm{m}^{2} / \\ \sin ^{2} 2 \theta \end{gathered}$ | Thres. <br> $\mathrm{N}_{\text {hit }}$ | $\text { Year } 1$ES | Year $2+3$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | ES | CC | $\mathrm{B}+\mathrm{NC}$ |
| vac. | 40 | 0.91 (.07) | 1.16 (.10) | 1.15 (.11) | 0.70 (.22) |
|  | 50 | 0.93 (.07) | 1.21 (.10) | 1.01 (.11) | 0.6 (1.0) |
|  | 60 | 1.00 (.09) | 1.18 (.13) | 0.98 (.12) | 1.6 (7.2) |
|  | 70 | 1.00 (.12) | 0.89 (.18) | 0.81 (.15) | 75. (59.) |
| $\begin{gathered} 10^{-6} \\ .01 \end{gathered}$ | 40 | 0.95 (.05) | 1.09 (.07) | 1.05 (.09) | 0.85 (.25) |
|  | 50 | 0.96 (.05) | 1.15 (.08) | 1.07 (.09) | -0.5 (1.6) |
|  | 60 | 0.99 (.07) | 1.09 (.10) | 1.02 (.10) | -1.5 (8.9) |
|  | 70 | 0.99 (.09) | 1.17 (.12) | 1.22 (.12) | 133. (71.) |
| $\begin{gathered} 10^{-5} \\ 0.3 \end{gathered}$ | 40 | 0.74 (.30) | 0.96 (.31) | 1.42 (.43) | 0.87 (.15) |
|  | 50 | 1.00 (.20) | 1.11 (.27) | 1.06 (.30) | 0.87 (.63) |
|  | 60 | 1.05 (.24) | 0.93 (.33) | 0.63 (.31) | 4.6 (3.1) |
|  | 70 | 1.00 (.30) | 0.77 (.35) | 0.80 (.40) | 14. (23.) |
| $\begin{gathered} 10^{-4} \\ .01 \end{gathered}$ | 40 | 0.89 (.11) | 1.02 (.14) | 1.30 (.23) | 0.85 (.17) |
|  | 50 | 0.89 (.12) | 1.08 (.17) | 1.06 (.25) | 0.81 (.71) |
|  | 60 | 1.00 (.19) | 1.43 (.27) | 1.56 (.37) | -3.4 (2.7) |
|  | 70 | 1.00 (.31) | 0.89 (.39) | 1.20 (.59) | -3.6 (17.) |

Table 6: Ratios of fitted numbers to Monte Carlo input numbers for $\mathrm{r} \leq 600 \mathrm{~cm}$

| Threshold |  | ES | CC | NC |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $N_{\text {hit }}$ | $E(\mathrm{MeV})$ |  |  | $\mathrm{D}_{2} \mathrm{O}$ | ${ }^{3} \mathrm{He}$ |
|  |  |  |  |  |  |
| 40 | 4.9 | 1000 | 8683 | 285 | 1742 |
| 50 | 6.2 | 656 | 6.994 | 89 | 1742 |
| 60 | 7.4 | 406 | 4464 | 13.7 | $17 \pm 2$ |
| 70 | 8.6 | 232 | 2653 | 1.04 | 1742 |

Table 7: Expected rates based on multiyear distributions inside $r=600 \mathrm{~cm}$ for Standard Solar Model with $5.8 \times 10^{6} \nu_{e}$ per $\mathrm{cm}^{2}$ per sec. The ${ }^{3} \mathrm{He}$ counters are assumed to have an efficiency of $85 \%$.

| $\Delta \mathrm{m}^{2} /$ | Thres. | Year 1 | Year 2 +3 |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\sin ^{2} 2 \theta$ | $\mathrm{~N}_{\text {hit }}$ | ES | ES | CC | NC |
|  |  |  |  |  |  |
| vac. | 40 | $.31(.03)$ | $.36(.03)$ | $.39(.04)$ | $3.9(1.2)$ |
|  | 50 | $.31(.02)$ | $.38(.03)$ | $.34(.04)$ | $1.4(2.8)$ |
|  | 60 | $.30(.03)$ | $.36\left(.0 \frac{1}{1}\right)$ | $.33(.04)$ | $3.1(14)$. |
|  | 70 | $.32(.04)$ | $.27(.03)$ | $.28(.05)$ | $180 .(142)$. |
|  | 40 | $.55(.03)$ | $.55(.04)$ | $.53(.04)$ | $4.7(1.4)$ |
| $10^{-6}$ | 50 | $.55(.03)$ | $.56\left(.0 \frac{1}{4}\right)$ | $.55(.04)$ | $-1.0(3.4)$ |
| .01 | 60 | $.55(.04)$ | $.55(.05)$ | $.54(.05)$ | $-2.9(18)$. |
|  | 70 | $.57(.05)$ | $.60(.06)$ | $.66(.07)$ | $-321 .(175)$. |
|  |  |  |  |  |  |
| $10^{-5}$ | 40 | $.04(.02)$ | $.05(.02)$ | $.08(.03)$ | $4.8(0.8)$ |
| 0.3 | 60 | $.05(.01)$ | $.06(.02)$ | $.06(.02)$ | $1.8(1.3)$ |
|  | 70 | $.04(.01)$ | $.05(.02)$ | $.04(.02)$ | $9.0(6.1)$ |
|  |  |  | $.04(.02)$ | $.04(.02)$ | $33 .(56)$. |
| $10^{-4}$ | 40 | $.16(.02)$ | $.16(.02)$ | $.17(.03)$ | $4.7(1.0)$ |
| .01 | 50 | $.11(.02)$ | $.11(.02)$ | $.08(.02)$ | $1.7(1.5)$ |
|  | 60 | $.08(.01)$ | $.09(.02)$ | $.07(.02)$ | $-6.6(5.4)$ |
|  | 70 | $.06(.02)$ | $.03(.02)$ | $.03(.02)$ | $-8.7(40)$. |
|  |  |  |  |  |  |

Table 8: Ratios of fitted numbers to numbers expected from the Standard Solar Model with a $\nu_{e}$ flux of $5.8 \times 10^{6} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$ for $\mathrm{r} \leq 600 \mathrm{~cm}$


Figure la




 Fi! :


Arizu: nnit! osfrei




Homomatsu 8", 3.8 ns


Fio: tinnss
 Fyure 7a


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$7 c$

$7 d$



Figuse 8


Fyure 9


n: : Esnr.r.








